# A PIM-aided Kalman Filter for GPS Tomography of the Ionospheric Electron Content

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#### **Abstract**

We develop the formalism for a PIM-based functional for stochastic tomography with a Kalman filter, in which the inversion problem associated with four-dimensional ionospheric stochastic tomography is regularized. For consistency, GPS data is used to select dynamically the best PIM parameters, in a 3DVAR fashion. We demonstrate the ingestion of GPS (IGS and GPS/MET) data into a parameterized ionospheric model, used to select the set of parameters that minimize a suitable cost functional. The resulting PIM-fitted model is compared to direct 3D voxel tomography. We demonstrate the value of this method analyzing IGS and GPS/MET GPS data, and present our results in terms of a 4D model of the ionospheric electronic density.

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### 1 Introduction

In previous work [1, 2, 3], we analyzed GPS data to extract information about the ionospheric electron density distribution. We can think of this distribution as a field in space-time which we try to represent using the information provided by the data. Since the ionosphere produces delays in the phase and group propagation of radio waves, having an accurate description of the electron content in the ionosphere is essential to any endeavor that uses radio wave propagation (such as tracking and navigating). In this paper we describe a novel parameterized tomographic technique to perform ionospheric imaging using Global Positioning System signal delay information.

Climatological models of the ionosphere have existed for a while now, but it is only recently that they have been used to complement other sources of data, such as GPS, in the inversion process. For instance, one can use input from a climatological model such as PIM [5] to complement GPS data in the inversion process, and to compare the results to other data [4]. The parameters controlling the model are input directly, however, and are not estimated themselves. One could reason, however, that if the models were good enough they could used to infer these parameters given other sources of data, such as GPS ionospheric delay data. The resulting "best-fit" parameters should be related to the ones one can obtain by independent means.

Let us give a brief introduction to ionospheric tomography (more details can be found in [1, 2, 3]). Let  $\rho(r, \theta, \phi, t)$  be the function that describes the electron density in some region of space  $(r, \theta, \phi, \phi, t)$  are spherical coordinates) at some time t. We can rewrite it as

$$\rho(r,\theta,\phi,t) = \sum_{J} a_{J}(t) \,\Psi_{J}(r,\theta,\phi), \tag{1}$$

where the functions  $\Psi_J(r,\theta,\phi)$  can be any set of basis functions we like. The goal in the inverse problem is to find the coefficients  $a_J(t)$ . In the case of GPS ionospheric tomography we use the information provided by the GPS ionospheric delay data along the satellite-receiver rays  $l_i$  to obtain a set of equations,

$$y_i = \int_{l_i} dl \, \rho(r, \theta, \phi, t) = \sum_{J} a_J(t) \int_{l_i} dl \, \Psi_J(r, \theta, \phi), \tag{2}$$

one for each ray  $l_i$ . Here  $y_i$  is the observed quantity. This is a set of linear equations of the form Ax = y, where the components of the vector x are the unknown coefficients  $a_J(t)$ . Assume that some cut-off in the basis function expansion is used and, therefore, that the x-space is N-dimensional. Let the y-space be M-dimensional (M is thus the number of data points). Since this system of equations may not have a solution we seek to minimize the functional  $\chi^2(x)$ , where (assuming uncorrelated observations of equal variance)

$$\chi^{2}(x) = (y - Ax)^{T} \cdot (y - Ax). \tag{3}$$

In practice we find that although the number of equations is much greater than the number of unknowns, the unknowns, i.e., the array x, are not completely fixed by the data. A way to restrict the solution space is to add some a priori constraints to the problem, and this can be implemented using the Lagrange multiplier method. Here we propose using a climatological model, such as PIM, to fill the gaps in the data and "smooth" the solution. But in order to use a climatological model one must provide the necessary input parameters. It is certainly possible to use for these parameters values provided by experimental sources of data (e.g., the solar flux is a measurable quantity). As was mentioned above, such techniques have already been used [4]. But another way to proceed is to use the GPS data itself, together with the model, to fix these parameters. This is especially important if it is suspected that the model parameters are not truly physical. If nothing else, this is an interesting exercise that will test the validity of the model.

A climatological mode, such as PIM, maps the value of a set of parameters,  $\lambda_i$ , to the space  $\{x\}$ . Just as is done in variational weather modeling, we can picture minimizing the cost functional

$$J(\lambda_i) = \sum_{j} \left( O_j^{exp} - O[x(\lambda_i)]_j \right)^2, \tag{4}$$

where  $O_j^{exp}$  are the observables and  $O[x(\lambda_i)]_j$  the modeled observables, in our case the slant delays produced by the ionospheric electrons. If we think of the climatological model image as the space spanned by a set of empirical orthogonal functions (which is the case in PIM), we see that this approach is just as the one described before, in the sense that a finite basis set is used to fit the data and represent the solution. What a model like PIM does is to provide us with a set of empirically or theoretically optimized basis functions to represent the ionospheric electron content.

### 2 PIM-aided Kalman Filtering

Kalman filtering is a very useful technique when dealing with a dynamic process in which data is available at different times. It is a natural way to enforce smoothness under time evolution, and is especially useful in the case of ionospheric stochastic tomography, when the "holes" in the information that we have at a given time (because of the particular spatial distribution of the GPS constellation and the receptor grid) may be "plugged" by the data from previous and future measurements. Indeed, in a Kalman filter we use the information contained in a solution to the inversion problem to estimate the next solution in the iteration process. In the study of the ionosphere, for example, we break the continuous flow of satellite delay data into blocks of a few hours, and simply model ionospheric dynamics by a random walk [7]. We can then process the data a given point in the iteration by asking that, to some extent, the solution be similar to the one

in the previous iteration, depending on how much confidence we have in that previous solution, and on how much we expect the dynamics to have changed things from one solution to the next. Here we complement this step by using the previous solution in the iteration process to fit a PIM model to the data. In other words, if  $x_n$  and  $C_n$  are the solution and the covariance matrix at epoch n, we first determine a minimum squares PIM fit. Let A be the observation matrix (which we know how to compute, given a grid). Then we minimize the cost functional

$$J(\lambda) = \left(y - A \cdot x^{PIM}(\lambda)\right)^2,\tag{5}$$

and this will determine the PIM parameters  $\lambda^i$ , and the resulting image,  $x_n^{PIM}(\lambda)$  and covariance matrix for the voxel image,  $C_n^{PIM}$ . This matrix is related to the covariance matrix for the PIM parameters,

$$C^{-1} = \nabla_{\lambda} \nabla_{\lambda'} J, \tag{6}$$

and is given by

$$C_n^{PIM} = \left(\nabla_{\lambda} x^i(\lambda) \left(\nabla_{\lambda} \nabla_{\lambda'} J\right)^{-1} \nabla_{\lambda} x^j(\lambda')\right)^{-1}.$$
 (7)

We will not worry too much about it for now, since it may be hard to compute these PIM derivatives. We will instead use an *ad hoc* covariance matrix, with the property that it will fill the holes in the data without affecting too much the solution where the data already provides some information (as is done in [4]).

Since the extremization equation for this functional is not linear and we could not easily compute derivatives we have chosen to minimize this functional using the Powell algorithm (see [8], for example).

Now, at epoch n+1 we are to minimize

$$\mathcal{K}_{n+1} = \chi_{n+1}^2(x_{n+1}) + \left(x_{n+1} - x_n^{PIM}(\lambda)\right)^T \left(C_{n+1}^{PIM} + \delta^2\right)^{-1} \left(x_{n+1} - x_n^{PIM}(\lambda)\right) \tag{8}$$

with respect to  $x_{n+1}$ . The parameter  $\delta$  (which will in general be a diagonal  $N \times N$  matrix) models the random walk away from the previous solution, and if of the form  $\delta^2 = \alpha \cdot t$ . Minimization yields

$$x_{n+1} = \left[ S_{n+1} + \left( C_n^{PIM} + \delta^2 \right)^{-1} \right]^{-1} \left( A_{n+1}^T y_{n+1} + \left( C_n^{PIM} + \delta^2 \right)^{-1} x_n^{PIM} \right), \tag{9}$$

where  $S_n = A_n^T A_n$ , and  $C_n^{-1} = S_n + \left(C_{n-1}^{PIM} + \delta^2\right)^{-1}$ . This can be easily implemented in an algorithm.

## 3 Ingesting GPS data into PIM versus using regular tomography

Let us first summarize our goals:

- To demonstrate the ingestion of GPS (IGS and GPS/MET) data into a parameterized ionospheric model, and to select the set of parameters that minimize a suitable cost functional.
- To compare the model fit to direct 3D voxel tomography.
- To develop a PIM-based functional for stochastic tomography with a Kalman filter, in which the inversion problem associated with four-dimensional ionospheric stochastic tomography is regularized. For consistency, GPS data is used to select dynamically the best PIM parameters, in a 3DVAR fashion.

GPS observables consist essentially of the delays experienced by the dual frequency signals ( $f_1 = 1.57542$  GHz and  $f_2 = 1.22760$  GHz) transmitted from the GPS constellation (25 satellites) and received at GPS receivers around the world and in orbit. Let  $L_i$  be the measured total flight time in light-meters of a ray going from a given GPS satellite to a receiver at the frequency  $f_i$  (including instrumental biases), and  $I = \int_{ray} dl \, \rho(x)$  be the integrated electron density along the ray (in electrons per square meter). Then  $L_i$  is modeled by  $L_i = D - I \, \alpha/f_i^2 + \tilde{c}_{sat} + \tilde{c}_{rec}$ , where  $\alpha = 40.3 \, m^3/s^2$ , D is the length of the ray, and  $\tilde{c}_{sat}$  and  $\tilde{c}_{rec}$  are the instrumental biases. In the present case we are interested in the frequency dependent part of the delay:  $L = L_1 - L_2$  (in meters). This is the derived observable and is modeled by  $(\gamma = 1.05 \times 10^{-17} \, m^3) \, L = \gamma \, I + c_{sat} + c_{rec}$ , independent of D (see [2] for more details). For the purposes of PIM-fitting, the solutions for the bias constants from the previous iteration are used to "fix" the observables delays, so that only the electronic part of the delay remains. At this point we have not tried to estimate the bias constants within the PIM-fitting analysis, although this should be possible. See the Appendix A for details on our bias constant treatment.

GPS data has been collected from GPS/MET and a subset of the International GPS Service (IGS) Network, for the day of February 23rd of 1997. This particular day has been chosen because of A/S is known to have been off. Geomagnetic and solar activity indices (as distributed by the US National Geophysical Data Center) for that day indicate a mean  $K_p$  index of 2.3, and  $F_{10.7} = 73$ .

The raw data has been pre-processed in order to obtain the observables using the procedures described in [2]. To describe the ionosphere we use five geocentric spherical layers beginning at 50 km above the mean surface (6350 km) of the Earth and extending 1300 km. Each layer consists then of two hundred voxels of dimensions 18° in latitude, times 18° in longitude, times 150 km of height for the first 4 layers.

The unknowns here consist of the electron densities at each of these voxels, plus the unknowns corresponding to the transmitter and receiver constant delays. These are estimated and used to correct the data prior to PIM-fitting. For a particular block, a minimum was found at  $F_{10.7} = 52$ 

and  $K_p = 0$ . Thus, we see that these parameters should not be taken as physical quantities but just as parameters in the model. The PIM fit had a reasonable quality (40 cm standard deviation). Using the parameters estimated form observation ( $F_{10.7} = 73$  and  $K_p = 2.3$ ) yields a standard deviation of 45 cm (they are far from the minimum). This is expected, as it is known that PIM tends to overestimate TECs ( $Rob\ Daniell$ ,  $private\ communication$ ).

### 4 Summary, Conclusions

In this paper we have summarized our efforts to use climatological models in tomographic analysis of GPS data. This is a more natural thing to try than one may think at first. After all, climatological models such as PIM are essentially the result of performing Empirical Orthogonal Function analysis using empirical or theoretical data, and in a way this is exactly what one would like to do in tomography: the basis functions used to span the space of possible solutions should be adapted to the field one is trying to map. Basis sets such as wavelets are a step in this direction, but they are optimized to attack more general problems, where certain characteristics of the field one is studying are known. Here we can refine the basis set even more, given the theoretical and experimental knowledge that we already posses about the ionosphere. We have seen that the parameters in the model are not really physical, and we conclude that it is necessary to perform such parameter fits prior using the model estimates in the Kalman filter. Future efforts should be directed towards the development of more refined parameterized models. The ingestion of GPS data into this type of model has been demonstrated here.

### Appendix A

Here we show how to take out the constants from the analysis. Let x denote the array solution, in which the first n entries correspond to the voxel unknowns, and thereafter to the bias constants. Let us rewrite  $x = x_{vox} + x_c$ , where  $x_{vox}$  is an array with zeros after the nth entry, and  $x_c$  has zeros until after the nth entry. Now,

$$\chi^{2}(x) = (y - Ax)^{T} \cdot (y - Ax)$$

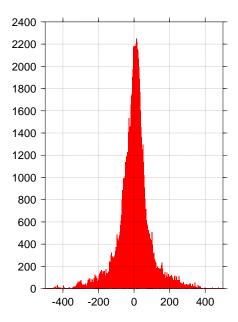
$$= y^{T}y + x_{c}^{T}A^{T}Ax_{c} - 2x_{c}^{T}A^{T}y + x_{vox}^{T}A^{T}Ax_{vox} + 2x_{vox}^{T}(-A^{T}y + A^{T}Ax_{c}).$$
(10)

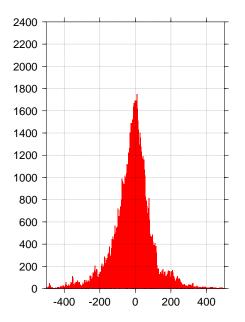
Hence, if we wish to fix  $x_c$ , all that is needed is to modify  $A^T y \to A^T (y - Ax_c) = (A^T y)_{corr}$ , and proceed without estimating the constants. Since  $x_{vox}$  is an array with zeros after the *n*th entry,

only the first n terms of  $(A^Ty)_{corr}$  are needed. The terms  $y^Ty + x_c^TA^TAx_c - 2x_c^TA^Ty$  are constants and do not affect the minimization solution. Hence we see that, up to irrelevant constant terms, the minimization problem is the same as without constants, but with a modified  $A^Ty$  term.

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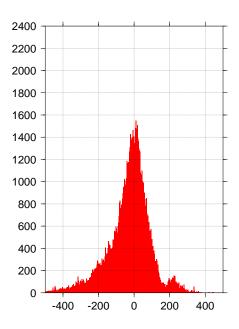


Figure 1: Left: Tomogrphic residual histogram. Standard deviation is 30 cm. Middle: PIM-fit residuals (at  $F_{10.7}=52$  and  $K_p=0$ ). Standard deviation is 40 cm. Right: PIM-fit residuals (at  $F_{10.7}=73$  and  $K_p=2.3$ ). Standard deviation is 45 cm.

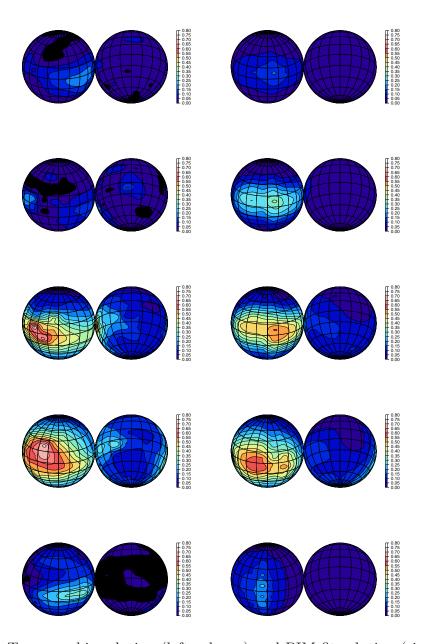


Figure 2: Tomographic solution (left column) and PIM-fit solution (right column), layer by layer and from bottom up, 6400-6550, 6550-6700, 6700-6850, 6850-7000, 7000-7700 km from center of Earth. Electronic density units are Tera electrons ( $10^{12}$ ) per cubic meter.